

I. Experimental test of non-macrorealistic cat states in the cloud by Huan-Yu Ku

Definition 1.1. Any macrorealistic system fulfill the **Leggett-Garg inequality** which is a result of the following two postulates: 1.) macrorealism and 2.) Noninvasive measurement. The former says that the a macroscopic system is in some definite state at any given time, and the latter states that some measurement can be performed to preserve the original state.

If a system exists in a definite state at any times, and the measurement is non-invasive, the probability of any outcomes, j , at times t_2 can then be written as the following equation, regardless of the system begin measured at time t_1 or not:

$$P_{t_2}(j) = P_{t_2}^M(j) = \sum_i P_{t_1}(i)P_{t_2}(j|i) \quad (1)$$

where the "invasiveness witness" is written as $W = |P_{t_2}(j) - P_{t_2}^M(j)|$, where W is 0 when the measurement is non-invasive

The noninvasive measurement is not realistic any may require some minor adjustment such that the state is changed when it's measured in time t_1 , the the equation (1) is modified to:

$$P_{t_2}(j) = P_{t_2}^M(j) = \sum_{i,k} \epsilon_{t_1}^M P_{t_1}(i)P_{t_2}(j|i) \quad (2)$$

then the "invasiveness witness" can be changed from the one with true non-invasiveness to the one with "clumsy non-invasiveness":

$$W = \max[|1 - \epsilon_{t_1}^M(i|i)|] \quad (3)$$

where it can be seen that if $\epsilon_{t_1}^M = 1$, representing that the existence of measurement does not change the weight of the outcomes probability at time t_1 , the outcomes probability changed to equation (1), where if the $\epsilon_{t_1}^M = 0$, the system is so strongly disturbed the state changed into some state that is orthogonal to the original state.

To quantify if the system is of macro-realistic or of quantum characteristic, the "disconnetivity" of a system is used.

Definition 1.2. disconnetivity is defined as the number of correlations between each "branch of super-position" one need to measure to distinguish the otherwise two indistinguishable states. The disconnetivity,

Γ can then be described as the following:

$$\Gamma = \max_i n_i \ni \frac{S_{n_i}}{\min_m (S_m + S_{n_i-m})} < \eta \quad (4)$$

Where S is the von-Neumann entropy and η is the some bound between classical mixtures and entangled states. ($\eta = 0.5$ is suggested for some reason unsepecified.)

To initialize a state with the most dsconnectivity(which is highly entangled), a cat state is generated:

Definition 1.3. A **cat state** is refered to a state that is in superposition of all-spin-up or all-spin-down states(GHZ state is one of them), the state can be generally written as:

$$|\phi\rangle_{t_1} = \cos\frac{\theta}{2} |0\rangle^{\otimes n} + \sin\frac{\theta}{2} |1\rangle^{\otimes n} \quad (5)$$

Assuming a non-invasive measurement is performed by a unitary and its conjugate. The "invasiveness witness" can be written as:

$$W = |P_{t_2}(j) - P_{t_2}^M(j)| = |P_{t_1}(i)P_{t_2}(j|i) - P_{t_1}(i)P_{t_2}^M(j|i)| = 1 - \cos^4\frac{\theta}{2} - \sin^4\frac{\theta}{2} \quad (6)$$

As a comparison, assume a state with no entanglement:

$$|\psi\rangle = \frac{1}{\sqrt{2^n}}(|0\rangle + |1\rangle)^{\otimes n} \quad (7)$$

The "invasiveness witness" is written as:

$$W = 1 - \frac{1}{2^n} \quad (8)$$

NOTE The "invasiveness witness" in equation (8) can be acted as a dimensionality witness since the effective dimension of a cat state is very low since it's only dominated by two state: all-spin-up or all-spin-down states.